On de Sitter Spacetime

Yoshihisa Kitazawa KEK Theory Center and Sokendai H. Kitamoto, R. Kojima

De Sitter Spacetime



FIG. 1 de Sitter space as a hyperboloid. Time goes up.— Right: Penrose diagram. Horizontal lines represent threespheres.

Cosmological horizon

$$\begin{split} \ell &= \sqrt{\frac{3}{\Lambda}}.\\ \frac{ds^2}{\ell^2} &= -d\tau^2 + \cosh^2\tau\,d\Omega_3^2\\ \frac{ds^2}{\ell^2} &= -V(r)\,dt^2 + \frac{1}{V(r)}dr^2 + r^2d\Omega_2^2, \end{split}$$

 $V(r) = 1 - r^2,$



FIG. 2 Past and future event horizon (diagonal lines). The static slicing covers the interior of the cosmological horizon (shaded).

Entropy and Temperature of Event Horizons

A. Black holes Bekenshtein Hawking

$$T_{
m hor} = rac{\kappa}{2\pi}.$$

 $\kappa = (4M)^{-1}$

the first law of thermodynamics.

$$rac{1}{T_{
m hor}} = rac{\partial S_{
m hor}}{\partial M},$$

 $S_{
m hor} = rac{A}{4}.$

B. de Sitter space

Gibbons Hawking

$$A_0 = 4\pi\ell^2,$$

 $\ell = \sqrt{rac{3}{\Lambda}}$
surface gravity $\kappa = 1/\ell$
 $T_{
m dS} = rac{1}{2\pi\ell},$
 $S_0 = rac{A_0}{4} = rac{3\pi}{\Lambda}.$

$$\begin{aligned} & \text{Swampland} \\ & |\nabla V| \geq \frac{c}{M_p} \cdot V , \\ & \min\left(\nabla_i \nabla_j V\right) \leq -\frac{c'}{M_p^2} \cdot V , \\ & \epsilon_V = \frac{M_{\text{Pl}}^2}{2} \left(\frac{V'}{V}\right)^2, \quad \eta_V = M_{\text{Pl}}^2 \left(\frac{V''}{V}\right), \\ & [\epsilon]: \epsilon_V \geq \frac{c^2}{2} \quad \text{or} \quad [\eta]: \eta_V \leq -c'. \end{aligned}$$

Cosmological Constant

2D gravity $g_{\mu
u} = e^{\phi} \hat{g}_{\mu
u},$



$$\frac{25-c}{48\pi}\nabla^2\phi - \Lambda e^\phi = 0,$$

$$e^{\phi_c} = \left(\frac{1}{-H\tau}\right)^2, \quad H^2 = \frac{24\pi}{c-25}\Lambda.$$

IR logarithmic effects



$$\langle \phi^2 \rangle \sim -\frac{12}{c-25} \phi_c(t) \Rightarrow \langle e^{\phi} \rangle \sim e^{(1-\frac{6}{c-25})\phi_c(t)}.$$

 $H_{\text{eff}}^2 \sim H^2 e^{-\frac{6}{c-25}\phi_c(t)} \sim H^2 a(t)^{-\frac{12}{c-25}}.$

Geometric Entropy

$$\begin{split} \delta S &= -\delta I_0 \\ &= \delta \phi (\frac{\epsilon}{2G} - \frac{25-c}{48\pi}) \mu^{\epsilon} [\int_M R \sqrt{\hat{g}} \ d^D x - 2 \int_{\partial M} K \ \sqrt{\hat{\gamma}} \ d^{D-1} x] \end{split}$$

 $\delta S/\delta \phi = (C-25)/12$

$$S = (C-25)/12 \phi(t) = (C-25)/6 Ht$$

Cosmological constant decreases as entropy increases

Infrared logarithmic effects to scalar perturbation

$$\epsilon = \frac{\kappa^2}{4} \left(\frac{\dot{\hat{\varphi}}}{H}\right)^2 = \frac{1}{\kappa^2} \left(\frac{V'(\hat{\varphi})}{V(\hat{\varphi})}\right)^2.$$

$$\int d^4x a^4(\tau) V'(\hat{\varphi}) \tilde{\varphi} \Big\{ 1 - \frac{11}{32} \frac{\kappa^2 H^2}{4\pi^2} \log a(\tau) \Big\}.$$

$$\frac{1}{\kappa^2} \frac{1}{V(\hat{\varphi})^2} \Big\{ 1 + \frac{9}{4} \frac{\kappa^2 H^2}{4\pi^2} \log a(\tau) \Big\}.$$

$$\left\{1 + \frac{25}{16} \frac{\kappa^2 H_*^2}{4\pi^2} \log a(\tau_*)\right\} \epsilon_*.$$

$$\eta = \frac{2V''(\hat{\varphi})}{\kappa^2 V(\hat{\varphi})}.$$

$$\int d^4x \ a^4 V''(\hat{\varphi}) \tilde{\varphi}^2 \{1 - \frac{3}{4} \frac{\kappa^2 H^2}{4\pi^2} \log a(\tau)\}.$$

$$\frac{1}{\kappa^2 V(\hat{\varphi})} \{1 + \frac{3}{4} \frac{\kappa^2 H^2}{4\pi^2} \log a(\tau)\}.$$

Cancelation of IR effects to this parameter!

Generation of linear potential due to IR effects

$$\frac{1}{2H}\frac{\dot{\epsilon}}{\epsilon} = -\eta + 2\epsilon + \frac{25}{32}\frac{\kappa^2 H^2}{4\pi^2}.$$
$$\epsilon \sim \epsilon_0 e^{\frac{25}{16}\frac{\kappa^2 H_0^2}{4\pi^2}H_0 t},$$

$$V = \frac{6H_0^2}{\kappa^2} (1 - \sqrt{\epsilon}\kappa\hat{\varphi}).$$

We must observe tensor modes soon



Conclusions

de Sitter spacetime may be a swampland in string theory

Cosmological constant decreases as entropy increases $A_0 = \frac{3\pi}{3\pi}$

$$S_0 = \frac{A_0}{4} = \frac{3\pi}{\Lambda}.$$

Evolusion of our universe suggests deep connection between cosmological constant and entropy