

Compact objects as the catalysts for vacuum decays

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$$c = \hbar = M_G^2 = 1/(8\pi G) = 1$$

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How stable is our vacuum ?

- **Compact objects as the catalysts for vacuum decays**

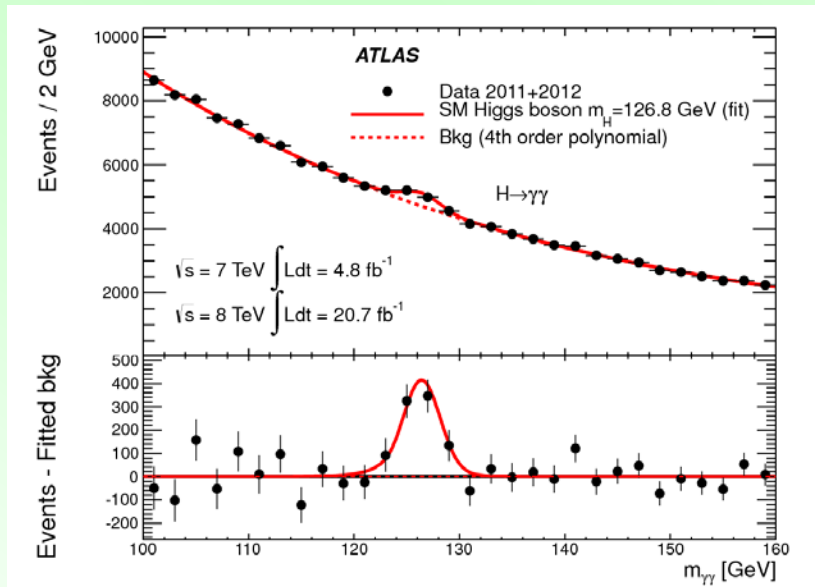
Formalism

Results

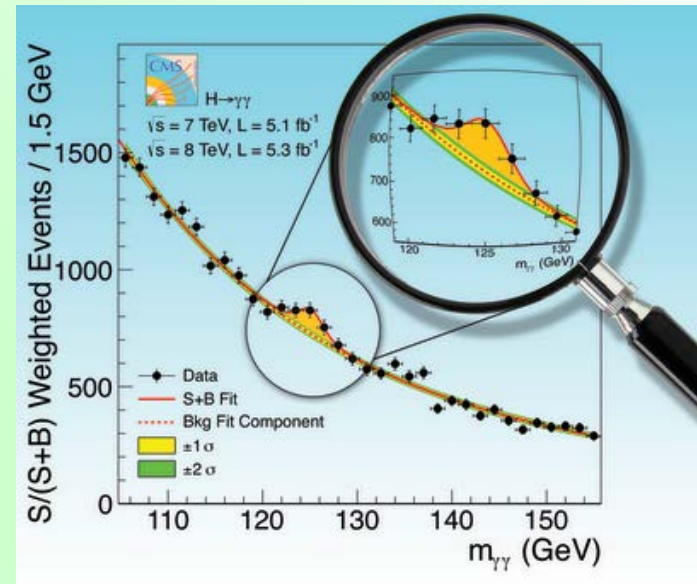
- **Discussion and conclusions**

Introduction

A Higgs particle was found !!



<http://www.atlas.ch/HiggsResources/>



<https://twiki.cern.ch/twiki/bin/view/CMSPublic/PhysicsResultsHIG>

A Higgs particle with 125 GeV mass has been finally found.

Running of the couplings

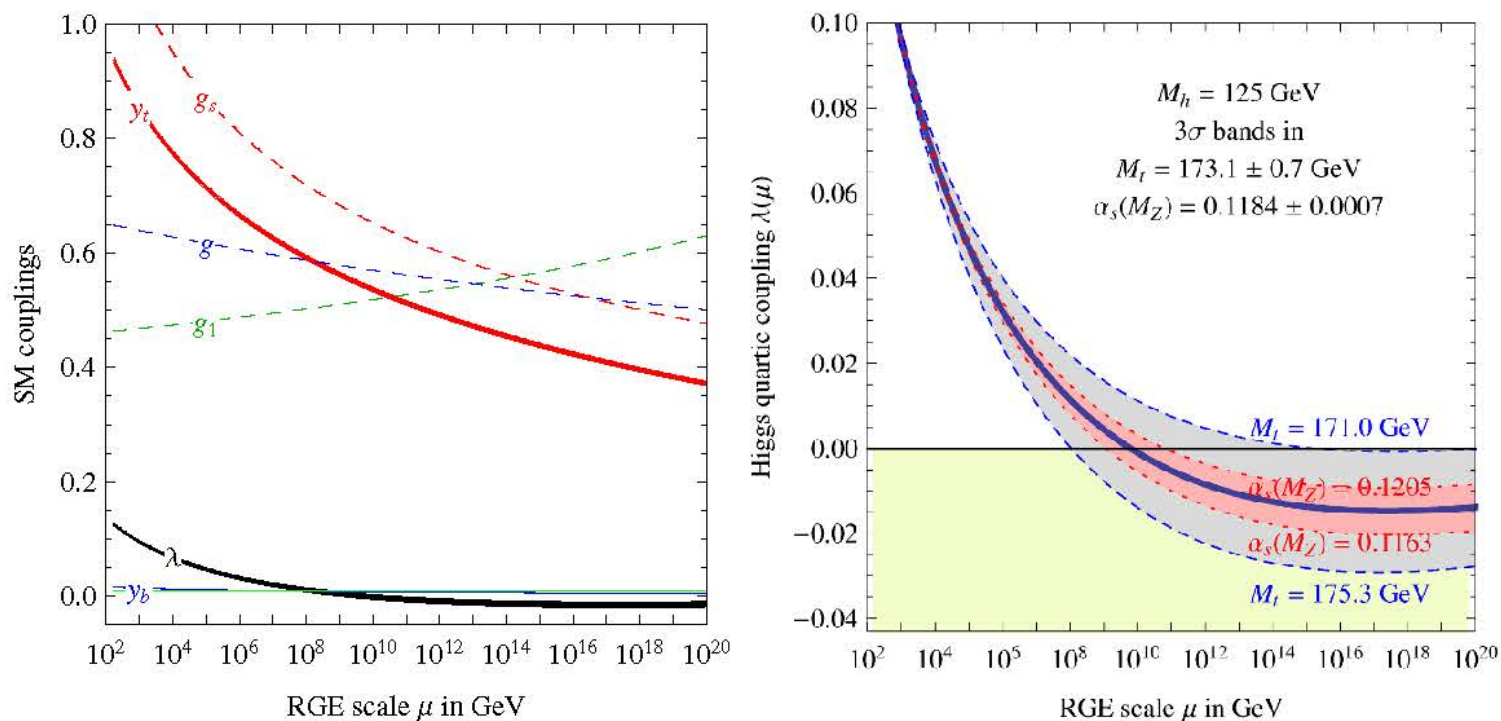


Figure 1. *Left:* SM RG evolution of the gauge couplings $g_1 = \sqrt{5/3}g'$, $g_2 = g$, $g_3 = g_s$, of the top and bottom Yukawa couplings (y_t, y_b), and of the Higgs quartic coupling λ . All couplings are defined in the $\overline{\text{MS}}$ scheme. The thickness indicates the $\pm 1\sigma$ uncertainty. *Right:* RG evolution of λ varying M_t and α_s by $\pm 3\sigma$.

(Degraasi et al. 2013)

The Higgs self-coupling can be negative at high energies.

Unstable vacuum ???

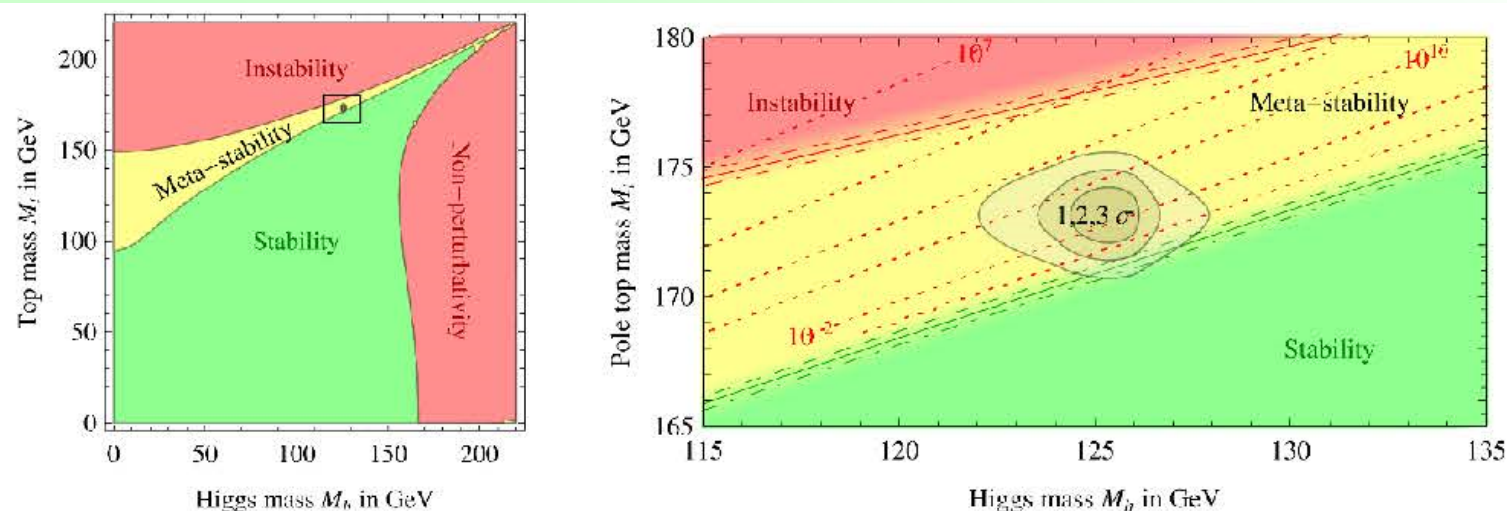


Figure 5. Regions of absolute stability, meta-stability and instability of the SM vacuum in the M_t - M_h plane. *Right:* zoom in the region of the preferred experimental range of M_h and M_t (the gray areas denote the allowed region at 1, 2, and 3 σ). The three boundaries lines correspond to $\alpha_s(M_Z) = 0.1184 \pm 0.0007$, and the grading of the colors indicates the size of the theoretical error. The dotted contour-lines show the instability scale Λ in GeV assuming $\alpha_s(M_Z) = 0.1184$.

(Degraasi et al. 2013)

Our vacuum might be in a meta-stable state, in which the time scale of the tunneling is longer than cosmic age,

Chigusa et al. shows that, with the best-fit values of the SM parameters, the decay rate of the EW vacuum per unit volume is about 10^{-554} /Gyr/Gpc³.

**These calculations are based on
the homogeneity and isotropy
of the Universe.**

(More precisely, $O(4)$ bounce)

**What happens if there is an impurity,
which can break homogeneity ?**

Black holes as bubble nucleation sites

Hiscock and, later, many people pointed out that (primordial) black holes can catalyze vacuum decays around them and significantly enhance the nucleation rate !!

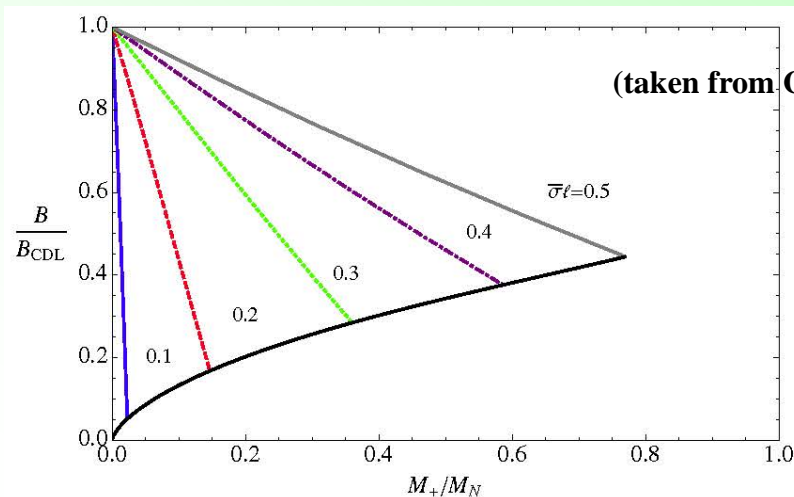
$$\Gamma \propto e^{-B} = e^{-B_0 + \Delta S}$$

↑
decay rate

↑
on-shell Euclidean action

↑
change of Bekenstein entropy < 0

(e.g. Gregory et al. 2014)



$B < B_{\text{CDL}}$ indicates enhancement of the decay rate.

**Is there any candidate for
such an impurity other than BHs ?**

Yes, there are many others !!!

**Compacts objects such as monopoles,
neutron stars, axion stars, oscillons,
Q-balls, black hole remnants,
gravastars, and so on...**

**Is there any crucial difference
between BHs and the other
compact objects ?**

Yes !!!

**While BHs have horizons,
the others not**

~~Black holes~~ as bubble nucleation sites

Compact objects

Hiscock and, later, many people pointed out that (primordial) black holes can catalyze vacuum decays around them and significantly enhance the nucleation rate !!

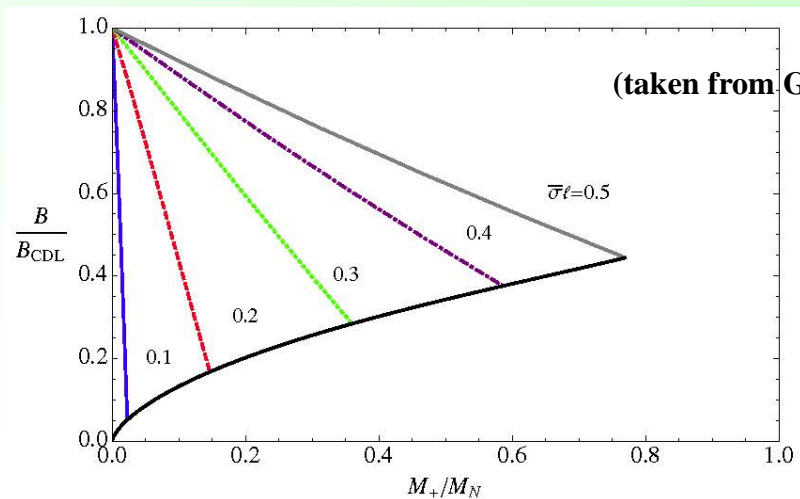
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↑
decay rate

↑
on-shell Euclidean action

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(e.g. Gregory et al. 2014)



$B < B_{\text{CDL}}$ indicates enhancement of the decay rate.

Formulation

Let us consider the nucleation of a thin wall vacuum bubble around a spherical object.

The metric inside and outside of the bubble :

$$ds^2 = -C_{\pm}(r_{\pm})dt_{\pm}^2 + D_{\pm}(r_{\pm})dr_{\pm}^2 + r_{\pm}^2 d\Omega_2^2$$

Specified by Einstein Eqs.

$$\left(C_{\pm} = D_{\pm}^{-1} = f_{\pm}(r_{\pm}) \equiv 1 - 2GM(r_{\pm})/r_{\pm} + H_{\pm}^2 r_{\pm}^2, \quad H_+ = 0, \quad H_-^2 \equiv -\frac{8\pi G}{3}\rho_v, \quad M(r_{\pm}) \equiv \int_0^{r_{\pm}} d\bar{r}_{\pm} 4\pi \bar{r}_{\pm}^2 \rho_c(\bar{r}_{\pm}) \right)$$

Israel's junction condition:

$$\left\{ \begin{array}{l} K_{AB}^{(+)} - K_{AB}^{(-)} = -8\pi G \left(S_{AB} - \frac{1}{2} h_{AB} S \right), \\ K_{AB}^{(\pm)} = \text{diag} \left(-\frac{d\beta_{\pm}}{dR}, \beta_{\pm} R, \beta_{\pm} R \sin^2 \theta \right), \\ S^A_B \equiv \text{diag} (-\sigma, p, p), \quad h_{AB} \equiv \text{diag} (-1, R^2, R^2 \sin^2 \theta), \end{array} \right.$$

Energy momentum tensor on the wall induced metric on the wall

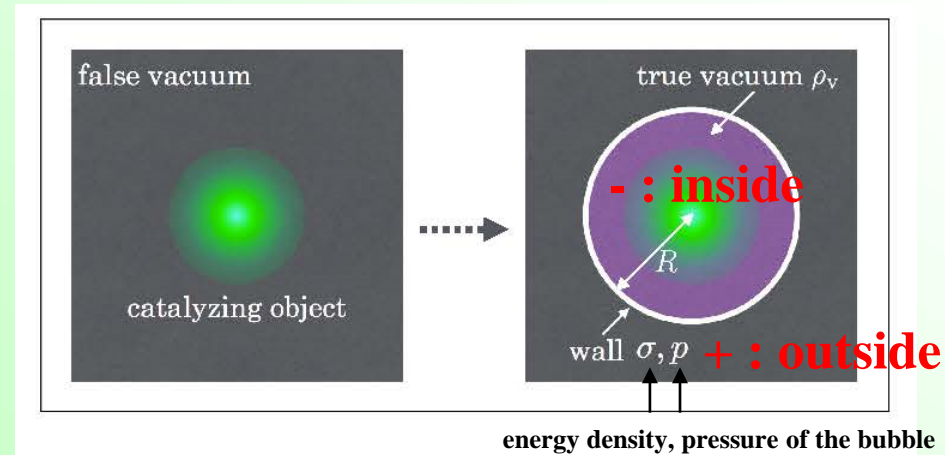


FIG. 1. A schematic picture showing a vacuum decay catalyzed by a static and spherical object.

($r_+ = r_- = R$ on the wall)

coordinate on the wall, $\xi^A = (\tau, \theta, \phi)$

$$\beta_{\pm} \equiv \epsilon_{\pm} \sqrt{f_{\pm} + (dR/d\tau)^2}$$

Formulation II

Israel's junction condition:

$$K_{AB}^{(+)} - K_{AB}^{(-)} = -8\pi G \left(S_{AB} - \frac{1}{2} h_{AB} S \right)$$

$(A=B=\tau)$
 $(A=B=0, \varphi)$

$$\left\{ \begin{array}{l} \frac{d}{dR} (\beta_- - \beta_+) = -8\pi G (\sigma/2 + p), \\ (\beta_- - \beta_+) = 4\pi G \sigma(R) R \\ (\beta_{\pm} \equiv \epsilon_{\pm} \sqrt{f_{\pm} + (dR/d\tau)^2}) \end{array} \right.$$

$$\sigma = m^{1-2w} R^{-2(1+w)},$$

(m : typical mass scale of wall, w = σ/p)

$$\left(\frac{dz}{d\tau'} \right)^2 + V(z) = -1,$$

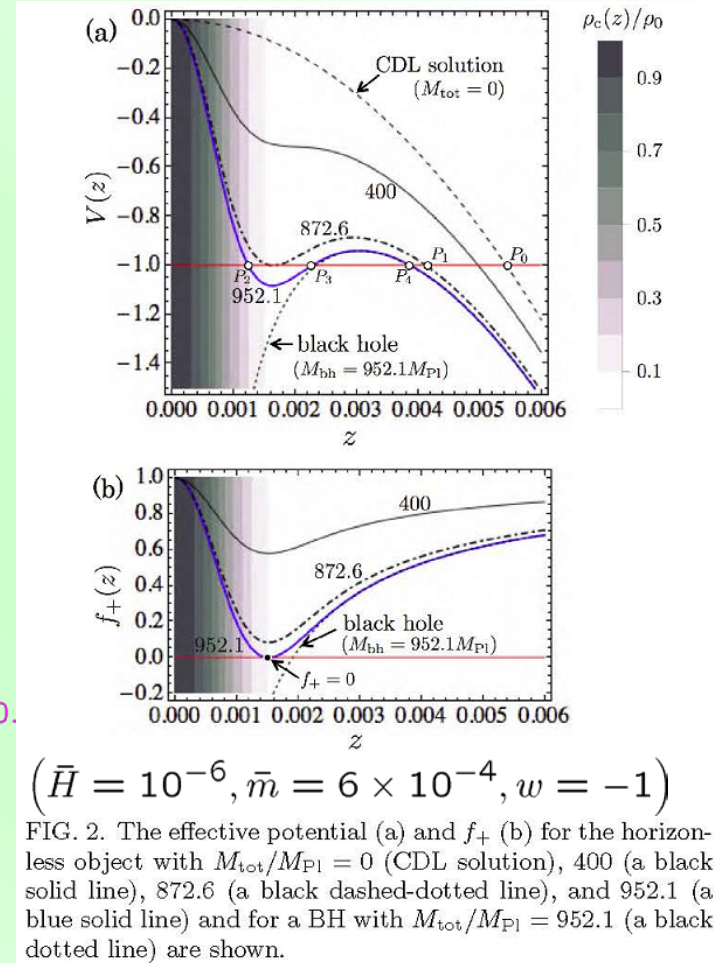
$$V(z) \equiv -\frac{a}{z} - \frac{z^2}{4} \left[\frac{z^{2(1+w)} \bar{m}^{2w-1}}{4\pi \bar{H}^{2w+1}} - \frac{4\pi \bar{H}^{2w+1}}{z^{2(1+w)} \bar{m}^{2w-1}} \right]^2 \leq 0.$$

$$z \equiv H_- R, \quad \tau' \equiv H_- \tau, \quad a \equiv 2GMH_-,$$

$$\bar{m} \equiv m/M_{\text{Pl}}, \quad \bar{H} \equiv H_-/M_{\text{Pl}}.$$

For the **nucleation (instanton)**, we need to take the **Euclidean** time.

Gaussian ansatz: $\rho_c(r) = \rho_0 e^{-r^2/\xi^2}$
(e.g. $\xi = 10^3 M_{\text{Pl}}^{-1}$)



$$M_{\text{tot}} \equiv \int_0^\infty dr' 4\pi r'^2 \rho_c(r') : \text{total mass of compact objects}$$

Decay rate

The Euclidean action B_{co} can be calculated from the bounce solution:

$$B_{\text{co}} = \frac{1}{4G} \oint d\tau_{\text{E}} (2R - 6GM + 2GM'R) \left(\frac{\beta_+}{f_+} - \frac{\beta_-}{f_-} \right).$$

Euclidean proper time of the wall

The decay rate :

$$\Gamma_{\text{D}} \sim R_{\text{CDL}}^{-1} \sqrt{\frac{B_{\text{co}}}{2\pi}} \exp(-B_{\text{co}}).$$



compare

$$\Gamma_{\text{C}} \equiv H_{\text{C}} \simeq 10^{-61} M_{\text{Pl}}$$

(the inverse of cosmological time)

Compactness parameter :

$$c \equiv \frac{\xi}{2GM_{\text{tot}}}. \quad \left(\rho_{\text{c}}(r) = \rho_0 e^{-r^2/\xi^2} \right)$$

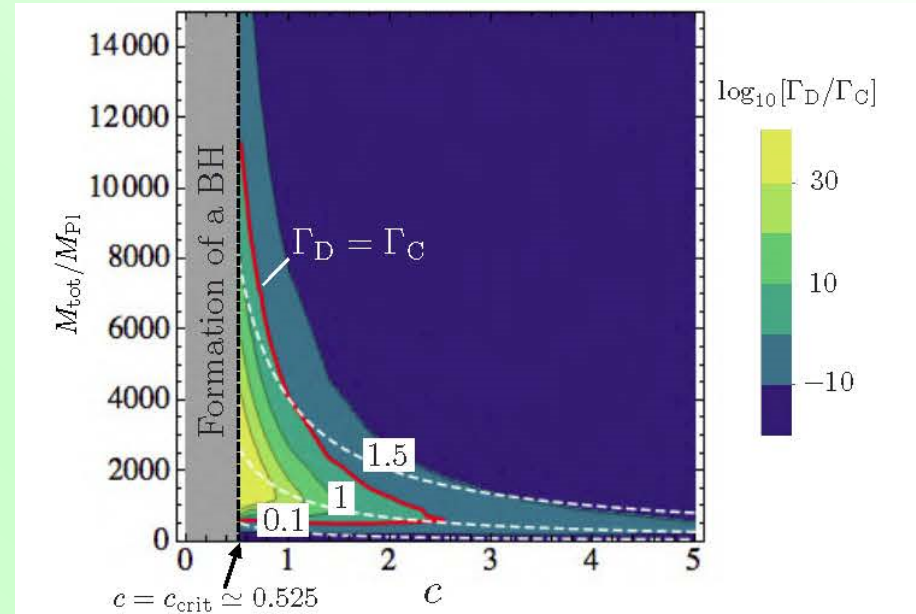


FIG. 3. A plot of the ratio of the decay rate, Γ_{D} , to the inverse of the cosmological time, Γ_{C} , as a function of the mass and compactness of the horizonless object. The contour of $\Gamma_{\text{D}} = \Gamma_{\text{C}}$ (red solid line) and contours of ξ/R_{CDL} (white dashed lines) are marked for reference. In the case of $c \leq c_{\text{crit}} \simeq 0.525$ (gray shaded region), the object inevitably collapses to a BH since a function $f_+(r)$ has zero points there.

The decay rate is drastically enhanced if the radius of the compact object ξ is comparable with the radius of CDL bubble and c is of the order of unity.

Comparison with the catalyzing effect of black holes

BH case:

$$\Gamma_D \sim R_{\text{CDL}}^{-1} \sqrt{\frac{I_E}{2\pi}} \exp(-B_{\text{bh}} + \Delta S).$$

$$(I_E = B_{\text{bh}} + |\Delta S|)$$

CO case:

$$\Gamma_D \sim R_{\text{CDL}}^{-1} \sqrt{\frac{B_{\text{co}}}{2\pi}} \exp(-B_{\text{co}}).$$

The decay rate with the **compact object** is **larger than** that with the **BH** thanks to **the absence of the decrement of Bekenstein entropy**.

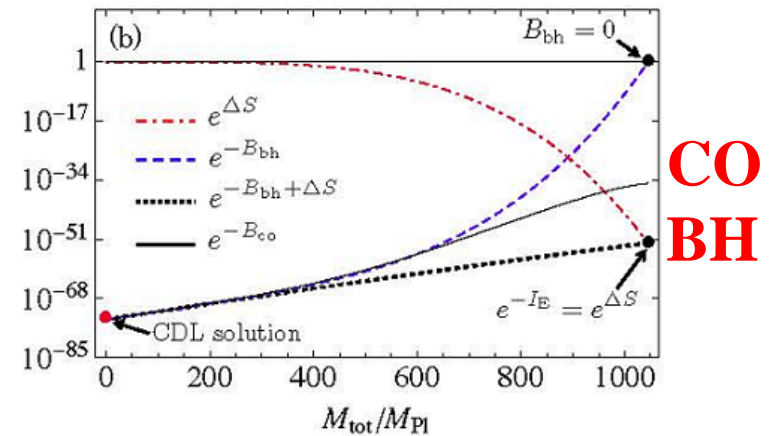
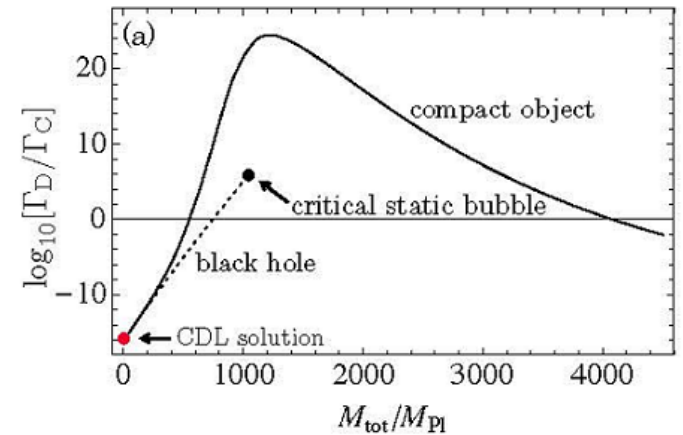


FIG. 4. The vacuum decay rates around a BH with mass M_{tot} (a dotted line) and that around a horizonless compact object, whose mass is M_{tot} and compactness is fixed with $\xi/2GM_{\text{tot}} = 1$, (a solid line) are shown. Red and black points show the decay rate of the CDL solution and a critical static solution, respectively.

Summary

- We have proposed a new possibility that **compact objects catalyze the vacuum decay efficiently**, which can give **yet another new constraint** on the abundance of compact objects.
- The bubble nucleation **rate is drastically enhanced around a compact object** if **the size of the horizonless object is comparable with the radius of CDL bubble and its compactness is of the order of unity**.
- We have shown that **the decay rate with the compact object is larger than that with the BH thanks to the absence of the decrement of Bekenstein entropy**.
- We have more tasks to do...