# Primordial black hole binaries as a probe of Hubble parameter

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#### 73.04 ± 1.04 km/s/Mpc



Years after the Big Bang



Image Credit: NAOJ



Image Credit: NAOJ

#### Primordial black holes as a potential candidate



Image Credit: ESA

#### How to observe the signals from PBHs?

Image Credit: ESO



$$h(t) = \frac{4}{d_L(z)} \left(\frac{G\mathcal{M}_Z}{c^2}\right)^{5/3} \left(\frac{\pi f(t)}{c}\right)^{2/3} \cos \Phi(t)$$
$$d_L(z) = \frac{1+z}{H_0} \int_0^z \frac{c}{E(z')} dz' \qquad \mathcal{M}_Z = \mathcal{M}(1+z)$$

Image Credit: ESO

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 $\Delta t = \frac{1}{H_0} \int_{z_1}^{z_2} \frac{dz}{E(z)(1+z)}$ 

Image Credit: ESO

How to construct redshift-distance (redshift-time) relation?

How to construct redshift-distance (redshift-time) relation? A statistical study on PBH binaries may help

## **PBH** formation



The primordial origin gives an identical primordial mass function

$$n(M) = \frac{f_{\text{PBH}}}{\sqrt{2\pi}\sigma M} \exp\left[-\frac{\ln^2(M/M_{pk})}{2\sigma^2}\right]$$

### **PBH** binary formation

The equation of proper separation *r* of two nearby PBHs with mass M is  $\ddot{r} - (\dot{H} + H^2)r + \frac{2M}{r^2}\frac{r}{|r|} = 0$ 



Ali-Haïmoud, Yacine, Ely D. Kovetz, and Marc Kamionkowski. "Merger rate of primordial black-hole binaries." *Physical Review D* 96.12 (2017): 123523.

### Primordial Black Hole Binaries as A Standard Timer

The initial probability distribution on a and e

$$\frac{dP}{dade} = \frac{3}{4} f_{\text{PBH}}^{3/2} \frac{a^{1/2}}{\bar{x}^{3/2}} \frac{e}{(1-e^2)^{3/2}}$$

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#### A single parameter standard timer

**Initial state:** Initial statistical distribution of dynamical systems dN

$$S(M;t_{\rm i}) = \frac{dN}{dM_{\rm i}}$$

Dynamics: time evolution of parameter in dynamical systems

$$\frac{dM}{dt} = -f(M) \Rightarrow \int_{M_{\rm i}}^{M_t} \frac{dM}{f(M)} = g(M_t) - g(M_{\rm i}) = -\Delta t$$

Later state: Statistical distribution of dynamical systems at physical time t

$$S(M;t) = \frac{dN}{dM_{\rm i}} \frac{dM_{\rm i}}{dM_t} = S(M;t_{\rm i}) \frac{g'(M_t)}{g'(g^{-1}(g(M_t) + \Delta t))}$$

**Observed state:** Redshifted statistical distribution of dynamical systems detected at redshift z

$$S_o(M_z; t) = S_o(M_z; t_i) \frac{g'(M_z)}{g'(g^{-1}(g(M_z) + \Delta t_z))}$$

**Redshift-time relation:** Comparing the observed state with the initial state gives the redshift-time relation

$$S_o(M_z;t) \simeq \begin{cases} S(M;t_i) \frac{dM_i}{dM_i(z)} &, \ g(M_z) \gg \Delta t_z \\ S_o(g^{-1}(\Delta t_z);t_i) \frac{g'(M_z)}{g'(g^{-1}(\Delta t_z))}, \ g(M_z) \ll \Delta t_z \end{cases}$$

#### A multi-parameter standard timer

Initial state: Initial statistical distribution of dynamical systems

$$S(\mathbf{M}; t_{i}) = \frac{dN}{d^{n}\mathbf{M}_{i}}$$

Dynamics: time evolution of parameter in dynamical systems

$$\frac{d\mathbf{M}}{dt} = -\mathbf{f}(\mathbf{M})$$

Later state: Statistical distribution of dynamical systems at physical time t

$$S(\mathbf{M};t) = \frac{dN}{d^{n}\mathbf{M}_{i}} \det \mathbf{J}(\mathbf{M},\Delta t) = S(\mathbf{M};t_{i}) \det \mathbf{J}(\mathbf{M},\Delta t)$$
$$\mathbf{J}_{ij} \equiv \partial M_{i}(t_{i}) / \partial M_{j}(t)$$

**Observed state:** Redshifted statistical distribution of dynamical systems detected at redshift *z* 

$$S_o(\mathbf{M}_z; t) = \frac{dN}{d^n \mathbf{M}_i(z)} \det \mathbf{J}(\mathbf{M}_z, \Delta t_z)$$

## How to extract the physical evolution time?

## The evolution of probability distribution in PBH binaries

$$\frac{dP}{da_t de_t} = \frac{dP}{da_i de_i} \det J(a, e, \Delta t)$$

$$J(a, e, \Delta t) = \begin{pmatrix} \partial a_i / \partial a_t & \partial a_i / \partial e_t \\ \partial e_i / \partial a_t & \partial e_i / \partial e_t \end{pmatrix}$$

$$\frac{da}{dt} = -\frac{128}{5} \frac{G^3 M_{\text{PBH}}^3}{c^5 a^3 (1 - e^2)^{7/2}} \left(1 + \frac{73}{24}e^2 + \frac{37}{96}e^4\right)$$

$$\frac{de}{dt} = -\frac{608}{15} \frac{G^3 M_{\text{PBH}}^3}{c^5 a^4} \frac{e}{(1-e^2)^{5/2}} \left(1 + \frac{121}{304} e^2\right)$$

## How to extract the redshift from the observable?



B. P. Abbott et al. Observation of Gravitational Waves from a Binary Black Hole Merger. Phys. Rev. Lett., 116(6):061102, 2016.

#### **Redshifted Chirp Mass**

$$\mathcal{M}_z = (1+z)\mathcal{M}$$





$$\Delta t = \int_{z_1}^{z_2} \frac{dz}{H(z)(1+z)}$$

$$H(z) = H_0 \sqrt{\Omega_{\gamma}(1+z)^4 + \Omega_m(1+z)^3 + \Omega_{\Lambda}}$$

$$\Delta t_2$$

$$\Delta t_2$$

$$\Delta t_2$$

## Merger rate of PBH binaries as a probe of Hubble parameter PBH mass function

$$n(M) = \frac{f_{\text{PBH}}}{\sqrt{2\pi}\sigma M} \exp\left[-\frac{\ln^2(M/M_{pk})}{2\sigma^2}\right]$$



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$$\mathcal{M}_{z} = (1+z)\mathcal{M}$$
$$d_{L} = \frac{1+z}{H_{0}} \int_{0}^{z} \frac{c}{E(z')} dz'$$

B. P. Abbott et al. Observation of Gravitational Waves from a Binary Black Hole Merger. Phys. Rev. Lett., 116(6):061102, 2016.

$$h(t) = \frac{4}{d_L(z)} \left(\frac{G\mathcal{M}_z}{c^2}\right)^{5/3} \left(\frac{\pi f(t)}{c}\right)^{2/3} \cos\Phi(t)$$



$$d_{L}^{i} = \frac{1 + z_{i}}{H_{0}} \int_{0}^{z_{i}} \frac{c}{E(z')} dz'$$

Assume a Hubble parameter  $\widetilde{H}_0$ 

$$z_i = z(d_L^i; \widetilde{H}_0)$$

$$m^i = m_z^i / (1 + z_i)$$





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detectable window function  

$$C(m_1^z, m_2^z) = \int_0^\infty \int_0^{\frac{m_1^z}{1+z}} \int_0^{\frac{m_2^z}{1+z}} n(m_1)n(m_2)W(m_1, m_2; z)p(z) dm_1 dm_2 dz$$
PBH mass function redshift distribution

$$P(m_1^z, m_2^z) = \frac{\partial^2 C(m_1^z, m_2^z)}{\partial m_1^z \partial m_2^z}$$

$$P(m_1^z, m_2^z) = \int_0^\infty n\left(\frac{m_1^z}{1+z}\right) n\left(\frac{m_2^z}{1+z}\right) W\left(\frac{m_1^z}{1+z}, \frac{m_2^z}{1+z}; z\right) \frac{p(z)}{(1+z)^2} dz$$

$$W(m_1, m_2; z) = \frac{N_{\text{obs}}(m_1, m_2; z)}{N_{\text{tot}}(m_1, m_2; z)} = \int_{a_{\min}}^{a_{\max}} \int_{e_{\min}}^{e_{\max}} P(a, e; z) \, dade$$

$$SNR = \sqrt{4 \int_{f_{\min}}^{f_{\max}} \frac{\left|\tilde{h}(f)\right|^2}{S_n(f)} df} > 8 \quad \tilde{h}(f) = \sqrt{\frac{5}{24} \frac{(G\mathcal{M}_z)^{5/6}}{\pi^{2/3} c^{3/2} d_L}} f^{-7/6}$$

$$p(z) \propto \frac{\dot{n}(z)}{1+z} \frac{dV_c}{dz}$$
  $\dot{n}(z) \propto \left(\frac{t(z)}{t_0}\right)^{-34/37}$ 



$$n(m) = \frac{1}{\sqrt{2\pi}\sigma m} \exp\left[-\frac{\ln^2(m/m_{pk})}{2\sigma^2}\right]$$

$$n(m) = \frac{\alpha - 1}{M} \left(\frac{m}{M}\right)^{-\alpha}$$

$$P_O(m_1^z, m_2^z) = \int_0^\infty n_p\left(\frac{m_1^z}{1+z}\right) n_p\left(\frac{m_2^z}{1+z}\right) W\left(\frac{m_1^z}{1+z}, \frac{m_2^z}{1+z}; z\right) \frac{p(z)}{(1+z)^2} dz$$

### **Gradient Descent Method**

$$P_O(m_1^z, m_2^z) = \int_0^\infty n_p\left(\frac{m_1^z}{1+z}\right) n_p\left(\frac{m_2^z}{1+z}\right) W\left(\frac{m_1^z}{1+z}, \frac{m_2^z}{1+z}; z\right) \frac{p(z)}{(1+z)^2} dz$$

$$P_T(m_1^z, m_2^z) = \int_0^\infty n' \left(\frac{m_1^z}{1+z}\right) n' \left(\frac{m_2^z}{1+z}\right) W\left(\frac{m_1^z}{1+z}, \frac{m_2^z}{1+z}; z\right) \frac{p(z)}{(1+z)^2} dz$$

$$E(n) = \sqrt{\frac{\sum_{1 \le i \le j, j=1}^{j=N} [P_T(m_1^z, m_2^z) - P_O(m_1^z, m_2^z)]^2}{(1+N)N/2}}$$

$$n_{k+1}(m_i) = n_k(m_i) - \gamma \frac{\partial E(n_k)}{\partial n_k(m_i)}$$



How about p(z)?

$$d_{L}^{i} = \frac{1 + z_{i}}{H_{0}} \int_{0}^{z_{i}} \frac{c}{E(z')} dz'$$
Assume a Hubble parameter  $\widetilde{H}_{0}$ 

$$z_{i} = z(d_{L}^{i}; \widetilde{H}_{0})$$

$$P_{O}(m_{1}^{z}, m_{2}^{z}) = \int_{0}^{\infty} \tilde{n}\left(\frac{m_{1}^{z}}{1+z}\right) \tilde{n}\left(\frac{m_{2}^{z}}{1+z}\right) W\left(\frac{m_{1}^{z}}{1+z}, \frac{m_{2}^{z}}{1+z}; z\right) \frac{p(z; \tilde{H}_{0})}{(1+z)^{2}} dz$$



However, we don't know the PBH mass function currently.

However, we don't know the PBH mass function currently. Another observable related with PBH mass function Merger rate of PBH binaries



$$R_{ij} = \rho_{\rm PBH} \min\left(\frac{n(m_i)}{m_i}, \frac{n(m_j)}{m_j}\right) \Delta_m \frac{dP}{dt}$$

$$\dot{N} = \int_{z_{\min}}^{z_{\max}} \frac{R}{1+z} \frac{dV_c}{dz} dz$$



Image Credit: NAOJ



Image Credit: NAOJ



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