
PHASE-ORDERING KINETICS IN FERROMAGNETIC SYSTEMS WITH LONG- RANGE INTERACTIONS

Seoul,
22/07/04

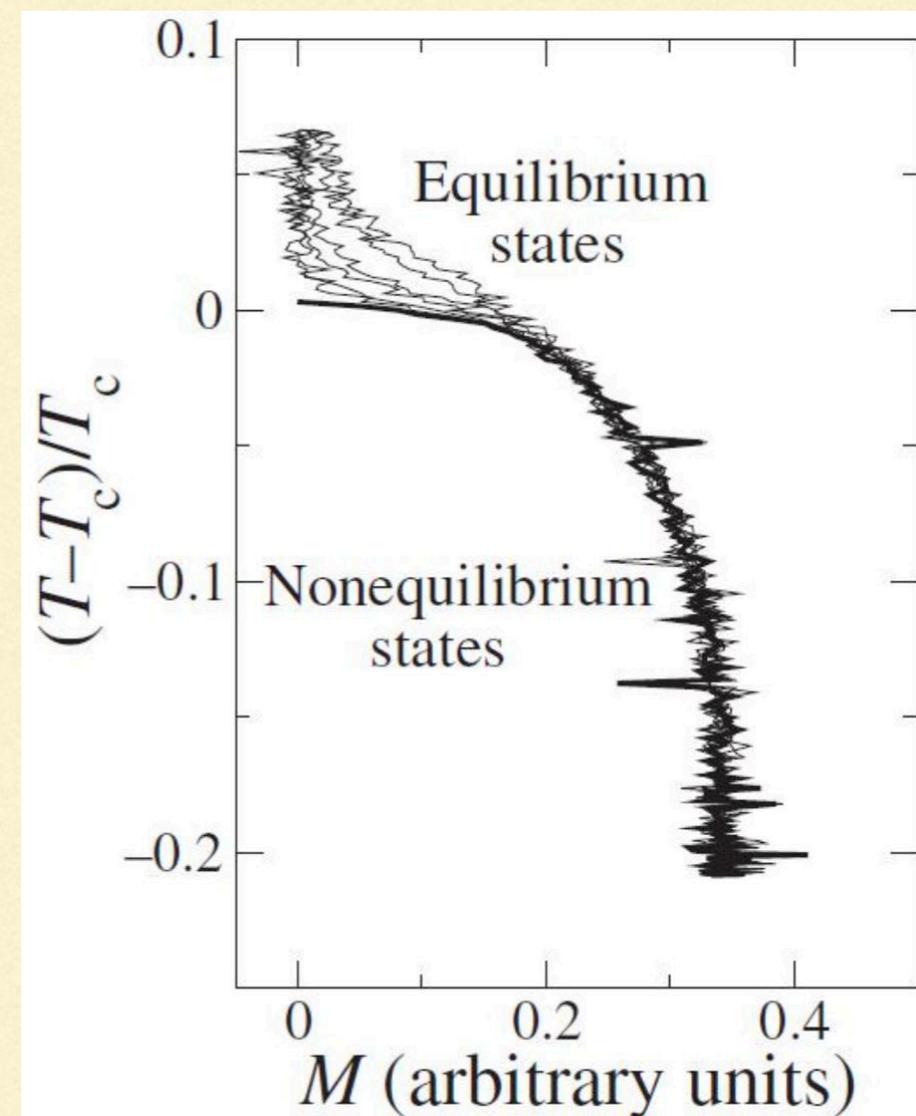
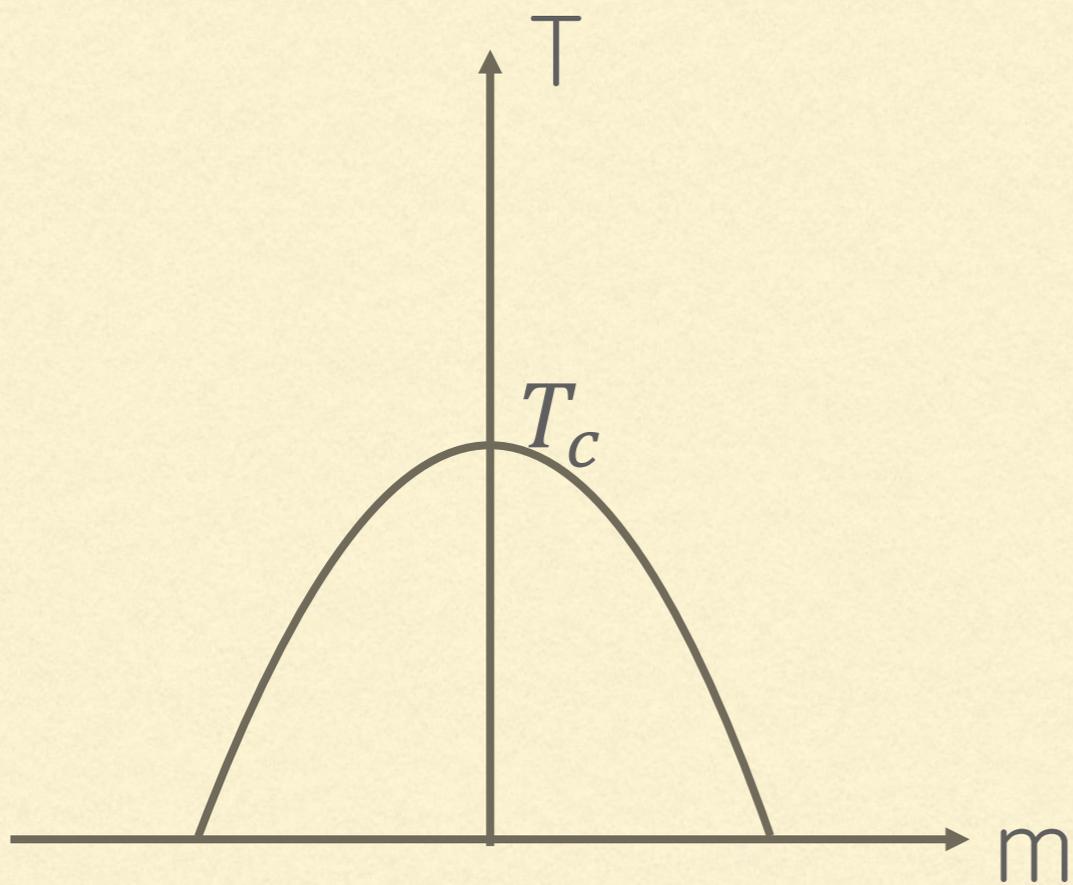
Work with:

audio Castellano (Institute for Complex Systems – CNR, Rom

Luca Smaldone (Salerno University)

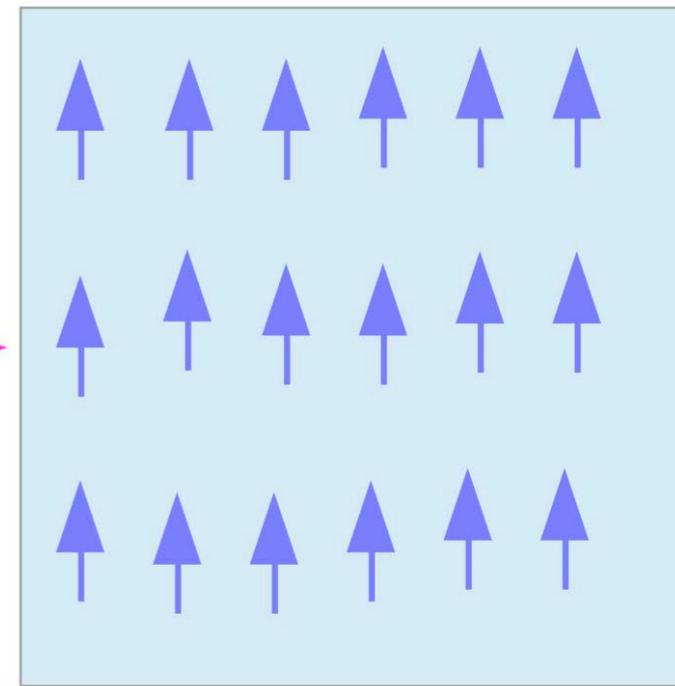
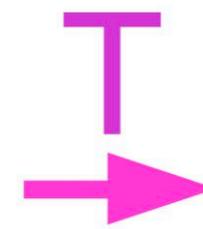
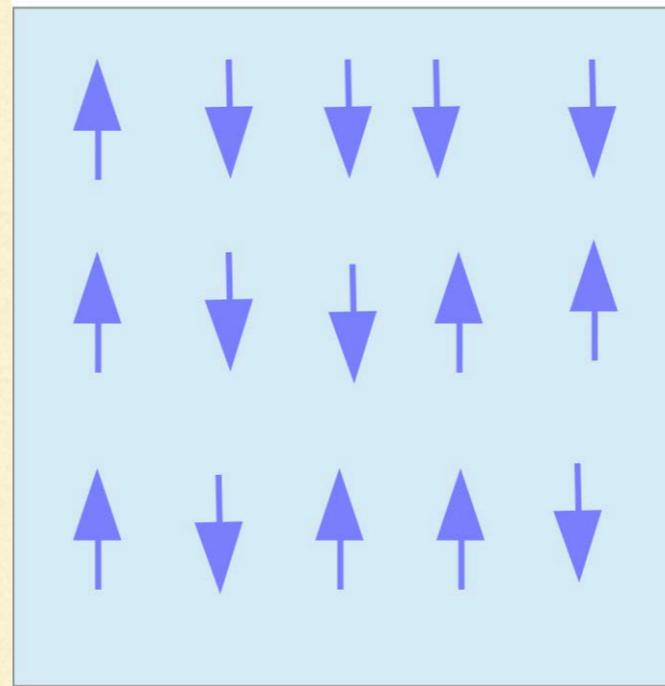
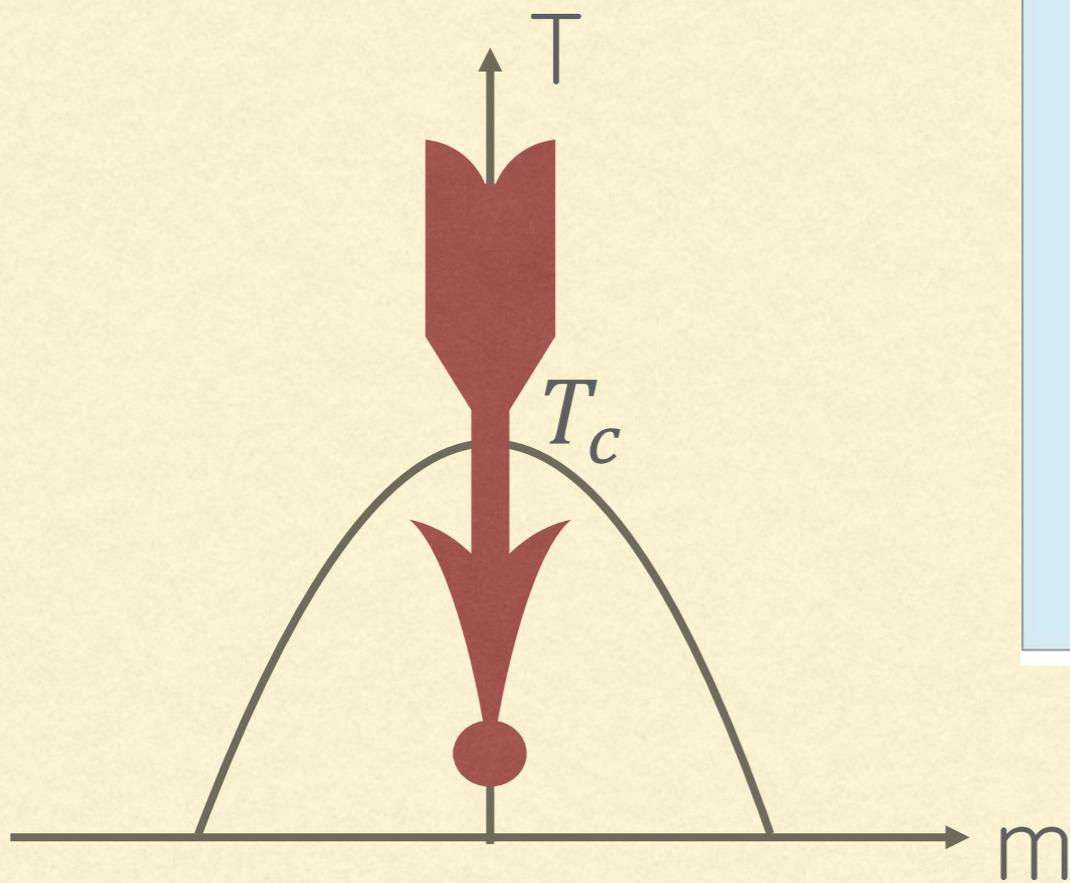
Salvatore dello Russo (Salerno University)

EQUILIBRIUM



Didascalia

NON-EQUILIBRIUM DYNAMICS (QUENCH)



Didascalìa

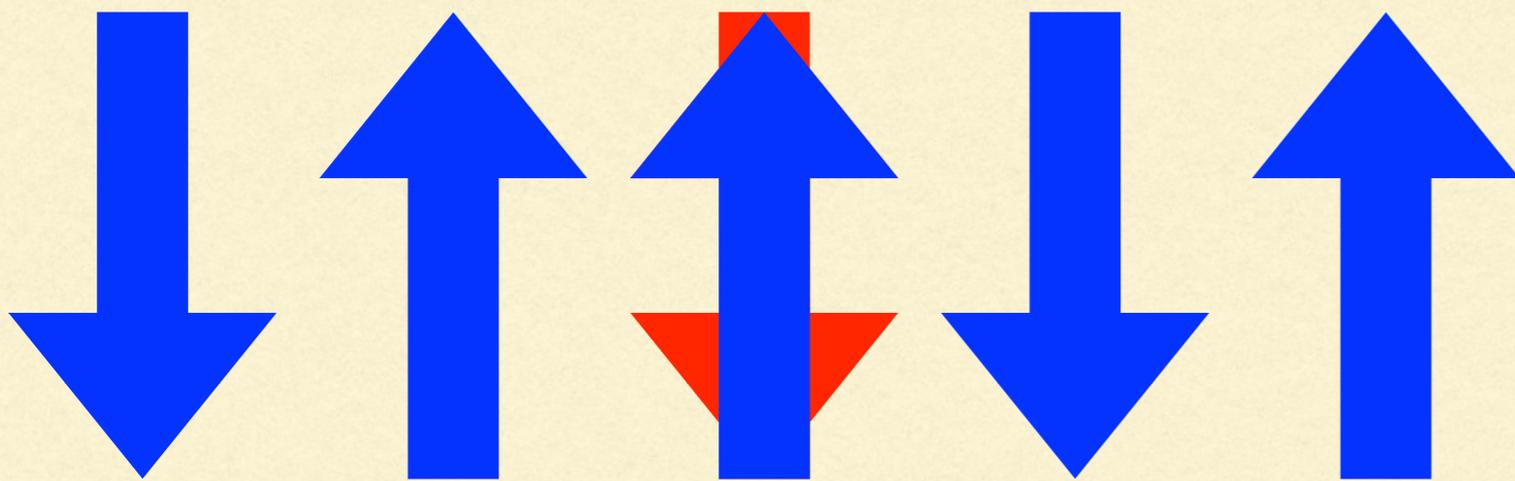
DISCRETE MODELS

ISING

$$-\mathcal{H}(\{s_i\}) = \sum_i \sum_r J(r) s_i s_{i+r}$$

$\omega_{\uparrow\downarrow}$

$$\frac{\omega_{\uparrow\downarrow}}{\omega_{\downarrow\uparrow}} = e^{-\beta\Delta E}$$

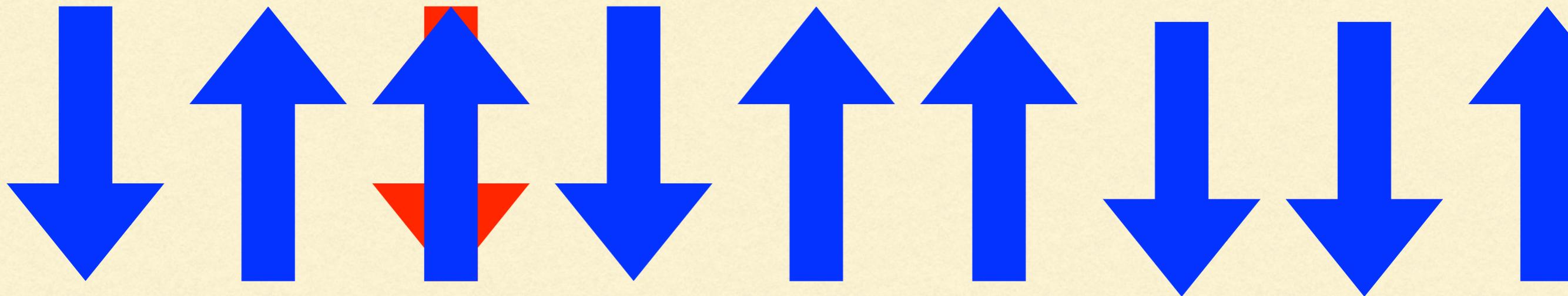


NCOP

DISCRETE MODELS

VOTER

$$-\mathcal{H}(\{S_i\}) = \sum_r J(r) s_i s_{i+r} \quad w(S_i) = \frac{1}{\omega_M} \sum_r P(r) \frac{\omega}{\omega_M} \sum_{k=i \pm r} (1 - S_i S_k)^{\beta \Delta E}$$

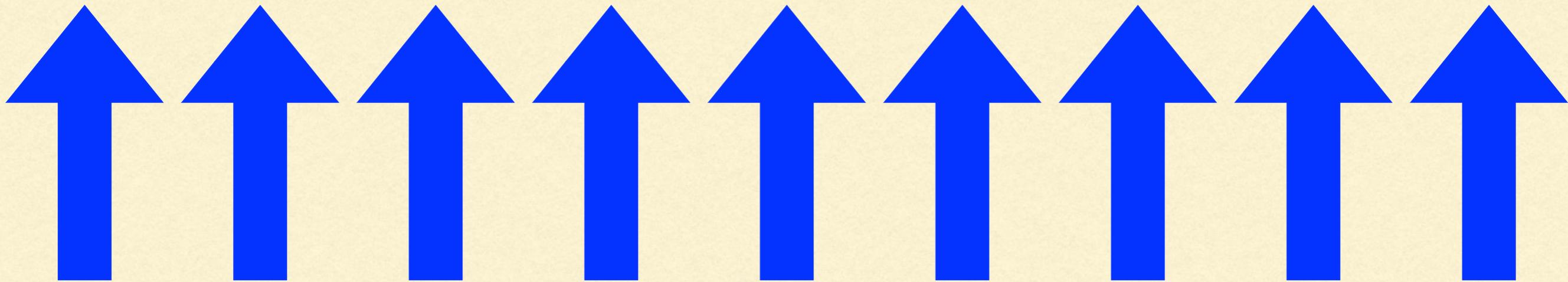


T=0

In D=1 with n.n. maps onto Ising $P(r) \sim J(r)$

Add Temperature

ABSORBING STATES (VOTER)

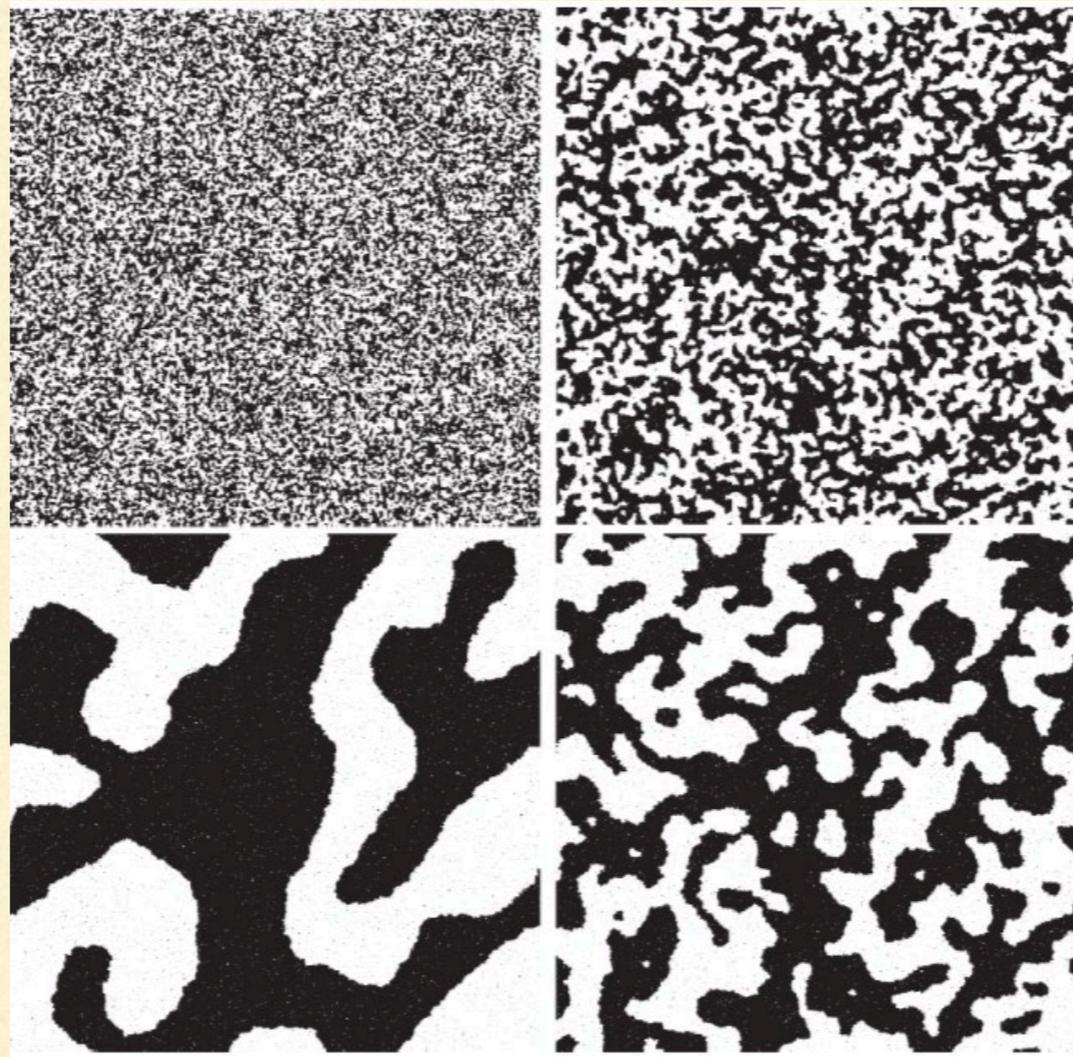


Fully ordered, consensus (up or down)

But, in some cases, there are stationary states without consensus,
with lifetime diverging with size

NEXT NEAREST NEIGHBOUR

ISING D=2



Time of
order
 N^z
to
equilibrate

$$L(t) \sim t^{1/2}$$

BEYOND NEAREST NEIGHBOURS (NEW RESULTS)

$$\mathcal{H}(\{s_i\}) = \sum_i \sum_{\vec{r}} J(r) s_i s_{i+\vec{r}}$$

$$J(r) = \begin{cases} J_0 e^{-r/R} \\ \frac{1}{Z} r^{-\alpha} \end{cases} \quad Z = \sum_r r^{-\alpha}$$

PHYSICAL REVIEW E

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Growth laws for phase ordering

- $\alpha = 0$, Mean field (MF) (our knowledge)
- $\alpha = \infty$, Nearest neighbours (NN)
- $\alpha > d$, rescaled asymptotics
- Only (well organised) continuum theory
- F. Corberi, E. Lippiello e P. Politi, *Phys. Rev. E* **119**, 26005 (2017).
Department of Theoretical Physics, The University of Manchester, M13 9PL, England
- F. Corberi, E. Lippiello e P. Politi, *Phys. Rev. E* **103**, 020102(R) (2020).
 (Received 5 March 1993)
- F. Corberi, E. Lippiello e P. Politi, *J. Stat. Mech.* 074002 (2019).
- F. Corberi, E. Lippiello e P. Politi, *Phys. Rev. E* **102**, 020102(R) (2020).
- R. Agrawal, F. Corberi, E. Lippiello, S. Puri e P. Politi, *Phys. Rev. E* **103**, 012108 (2021).
- F. Corberi, A. Iannone, M. Kumar, E. Lippiello e P. Politi, *SciPost Phys.* **10**, 109 (2021).
- F. Corberi, M. Kumar, E. Lippiello, e P. Politi, *Chaos, Solitons and Fractals* **173**, 113681 (2023).

BEYOND NEAREST NEIGHBOURS (NEW RESULTS, VOTER)

1D

F. C. and C. Castellano, J.Phys. Compl. 5, 025021 (2024)

F. C. and L. Smaldone, Phys. Rev. E 109, 034133 (2024)

F. C., S. dello Russo e Luca Smaldone, arxiv.org/abs/240

$$w(S_i) = \frac{1}{2N} \sum_r P(r) \sum_{k=i \pm r} (1 - S_i S_k)$$

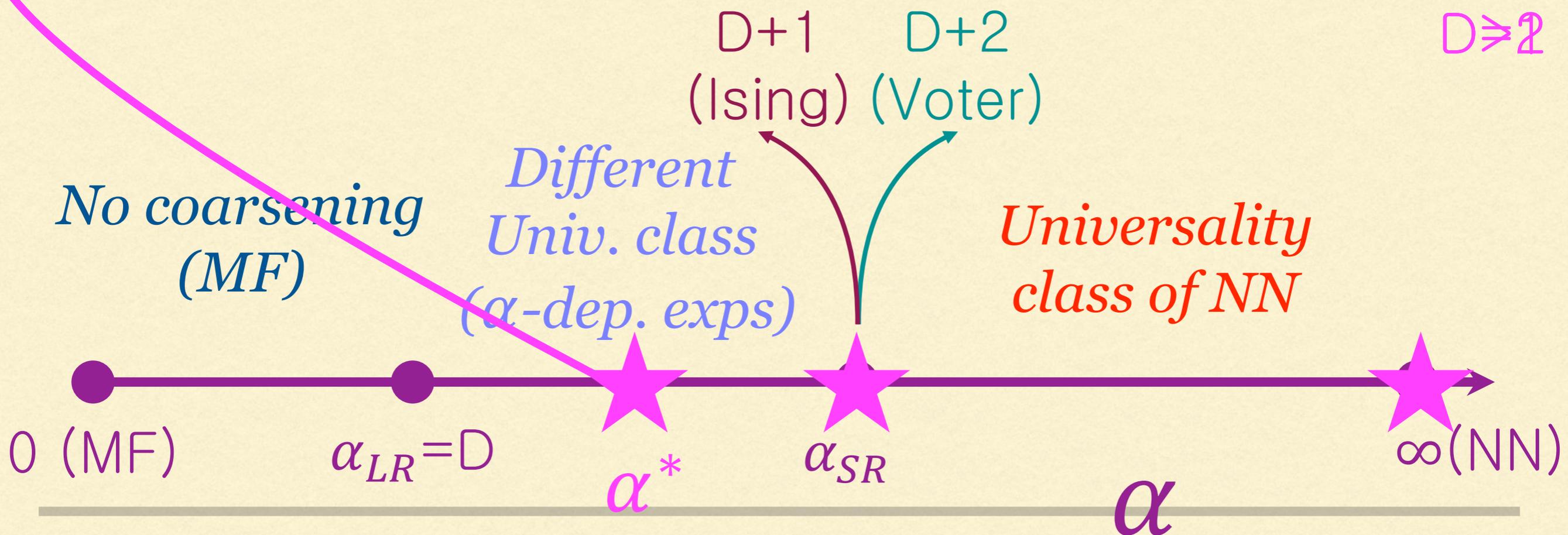
Solvable also for any D

$$\frac{d}{dt} \langle S_i S_j \rangle = -2 \langle S_i S_j [w(S_i) + w(S_j)] \rangle$$

$$\frac{d}{dt} C(r, t) = -2C(r, t) + 2 \sum_{\ell} P(\ell) [C(|r - \ell|, t) + C(r + \ell, t)]$$

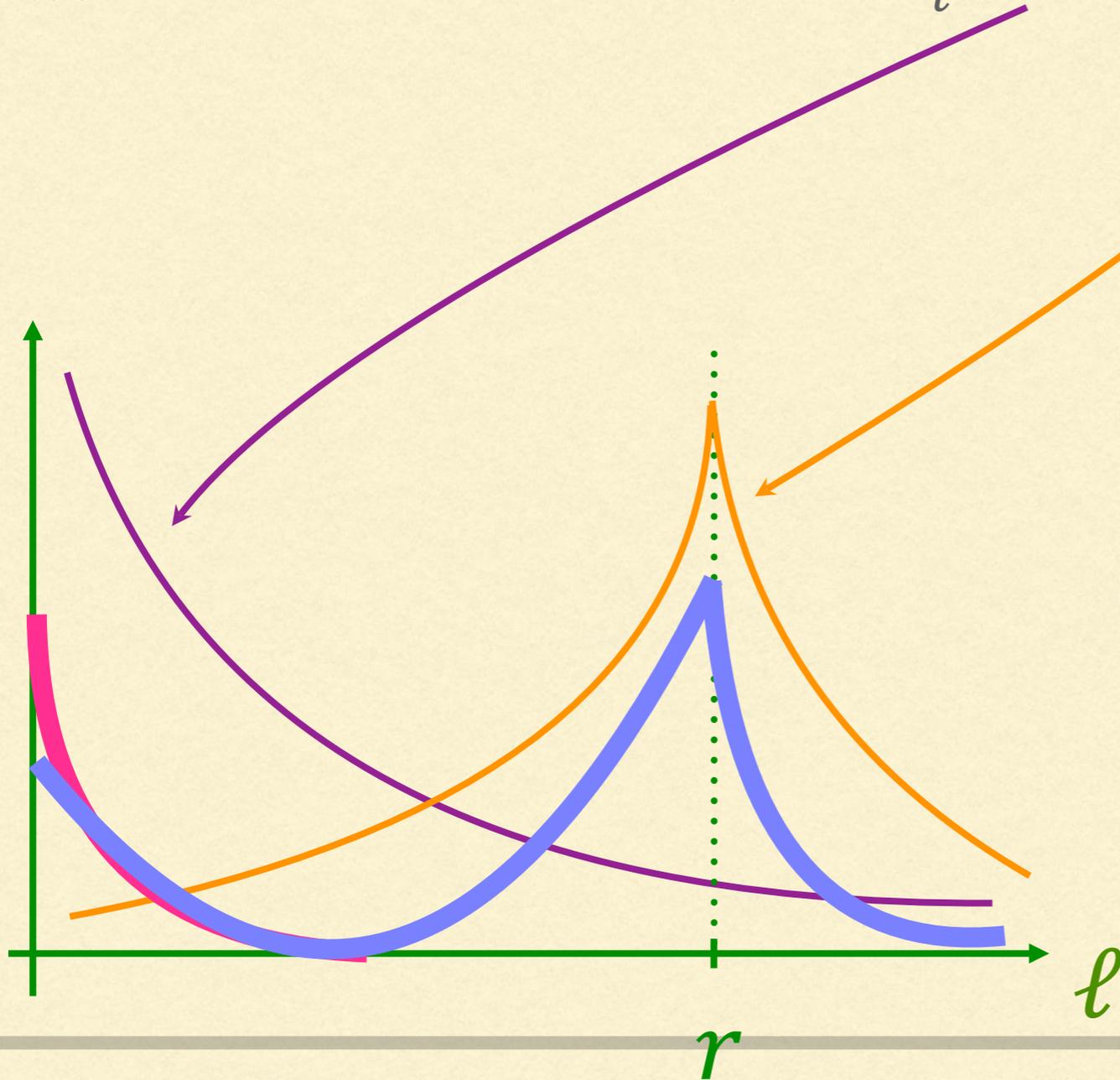
GENERAL COARSENING PROPERTIES

Only for voter: Highly correlated disordered stationary state
(non-integrable C, lifetime diverges with N)



VOTER

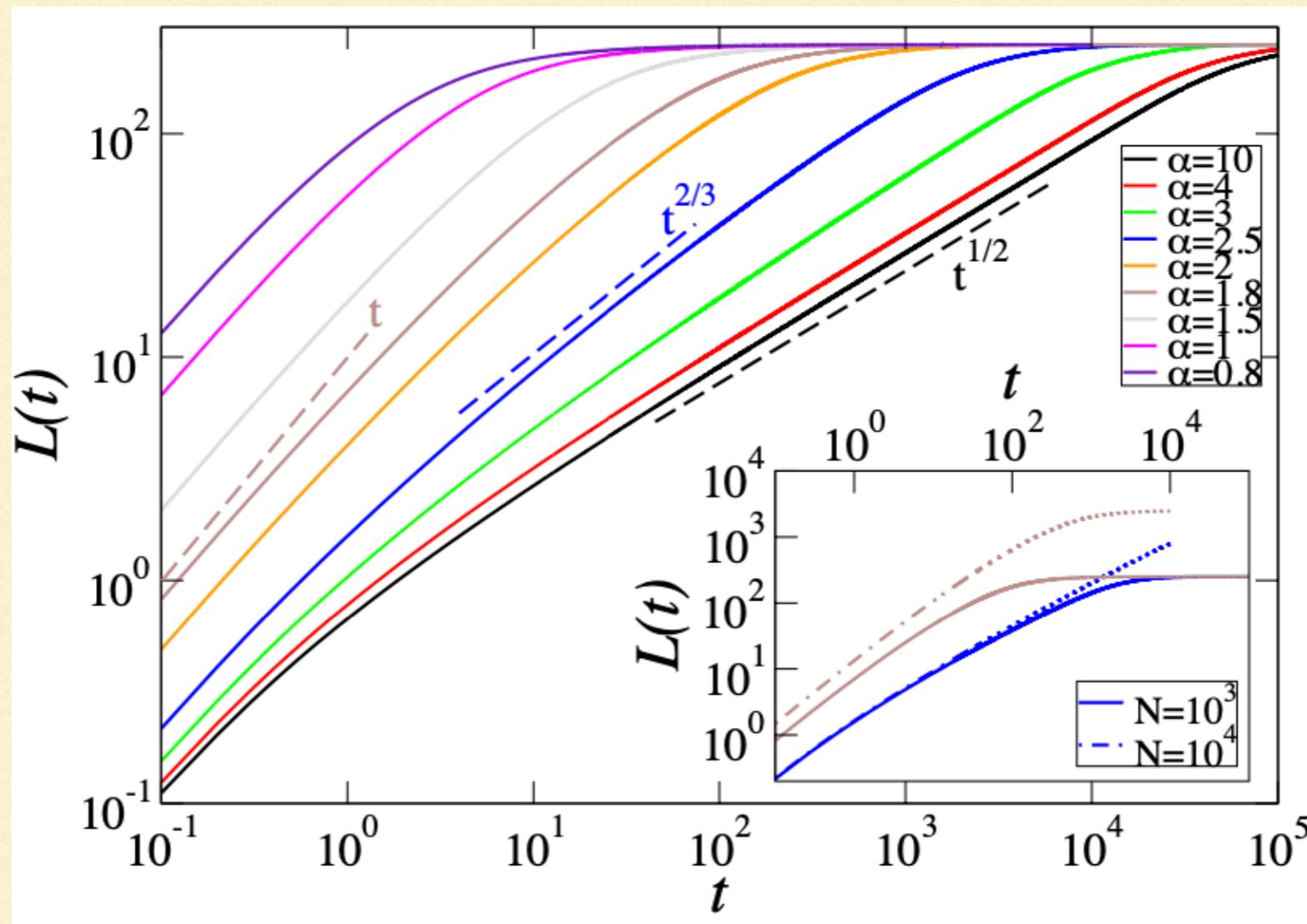
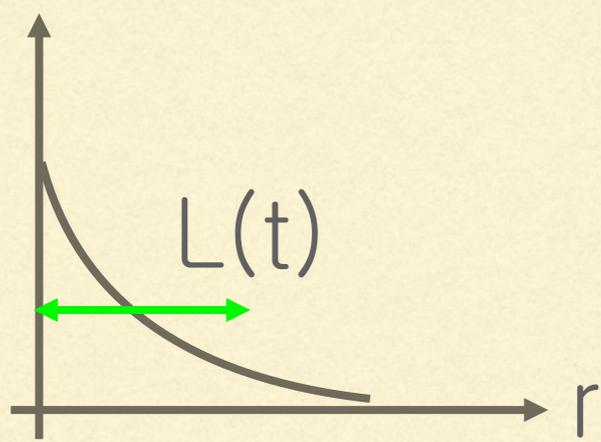
$$\frac{d}{dt} C(r, t) = -2C(r, t) + 2 \sum_{\ell} P(\ell) [C(|r - \ell|, t) + C(r + \ell, t)]$$



$$\alpha \leq \alpha_{SR}$$

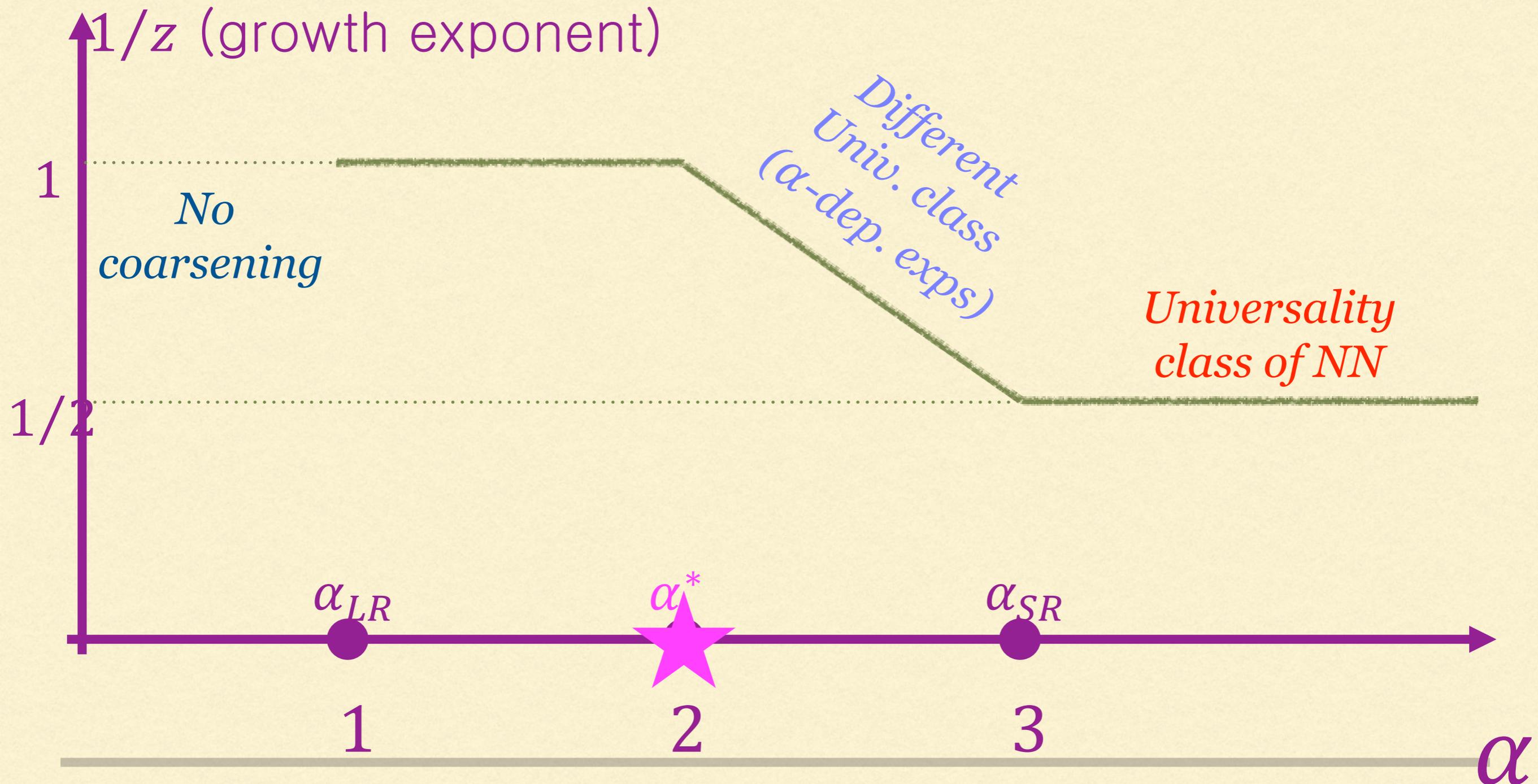
$C(r,t)$

VOTER



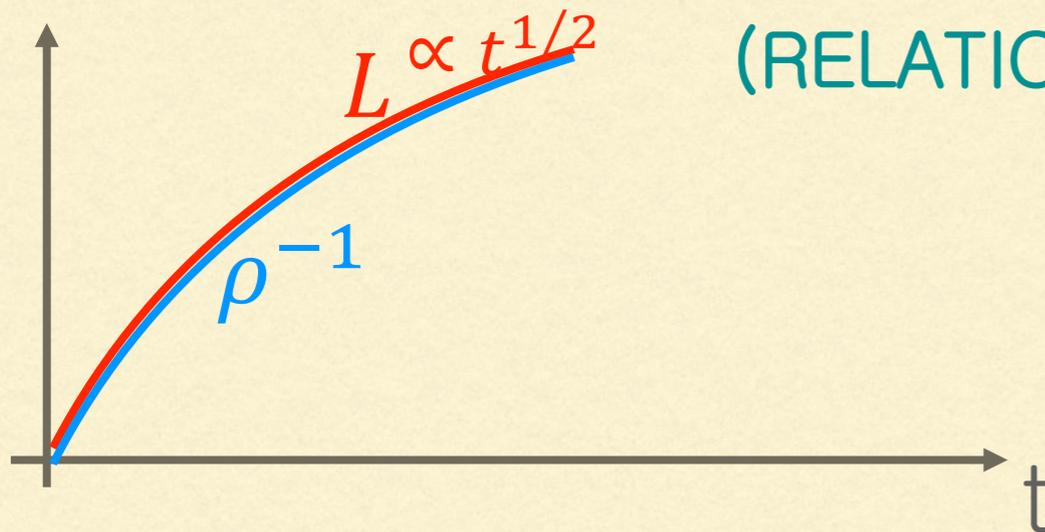
VOTER $D = 1$

SUMMARY OF GROWTH-LAW

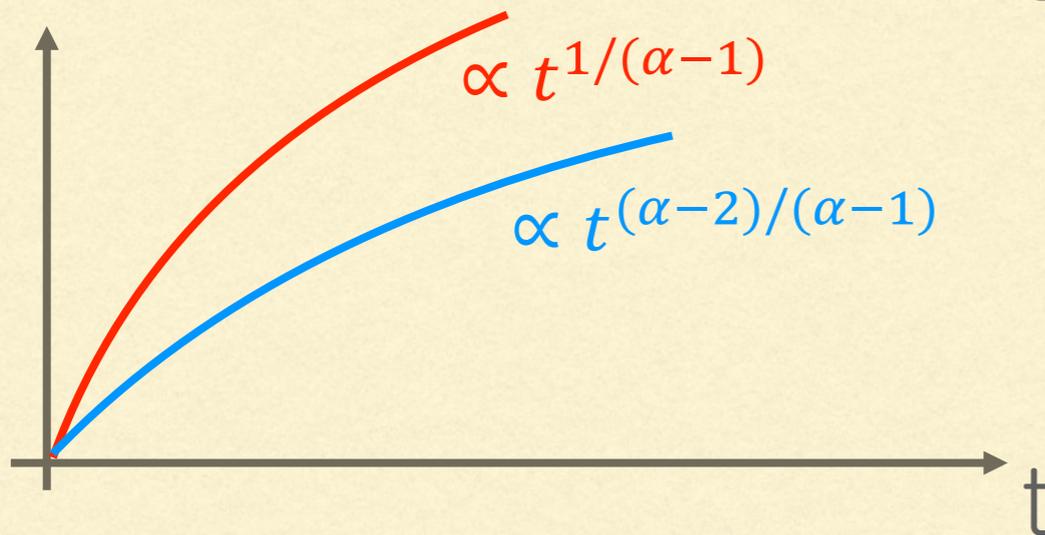


VOTER

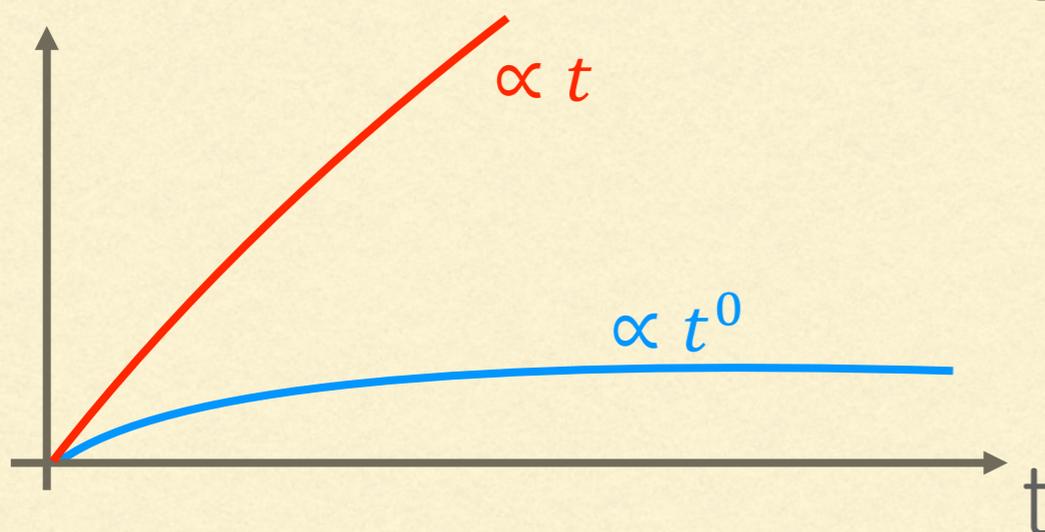
(RELATION $L(t) \leftrightarrow \rho^{-1}(t)$)



$$\alpha > \alpha_{SR}$$



$$2 < \alpha \leq \alpha_{SR}$$



$$\alpha < \alpha^* = 2$$



VOTER

(STATIONARY STATES)

$$C_{stat}(r) \propto r^{-(D-2)}, \text{ for } \alpha > \alpha_{SR} \text{ (for } D \geq 3)$$

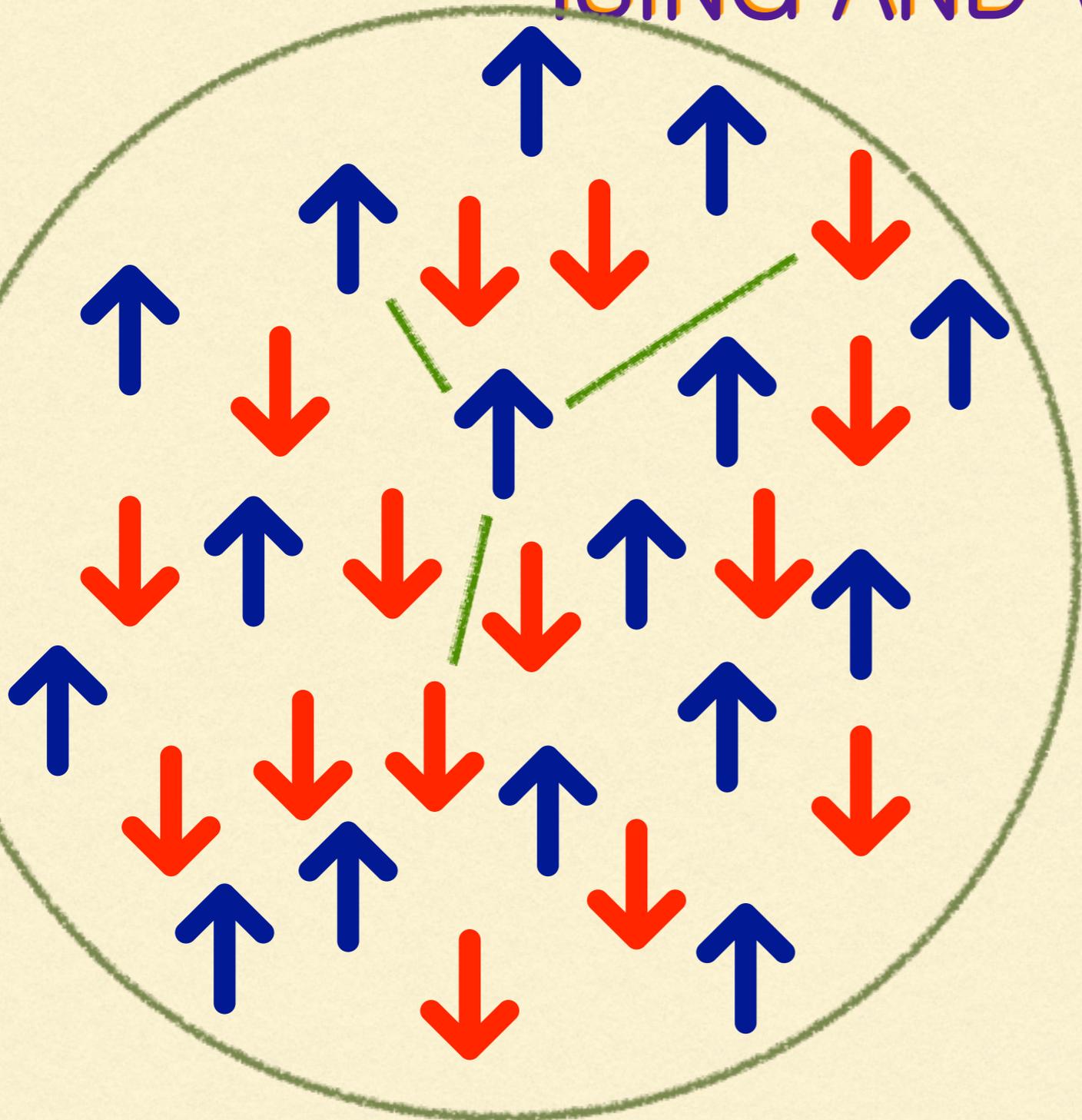
$$C_{stat}(r) \propto r^{-(2D-\alpha)}, \text{ for } \alpha_{LR} < \alpha \leq \min(\alpha_{SR}, \alpha^*)$$

$$C_{stat}(r) \propto r^{-\alpha}, \text{ for } \alpha \leq \alpha_{LR}$$

$$\lim_{N \rightarrow \infty} \int d^D r C_{stat}(r) = \infty \quad \Rightarrow \text{Strongly correlated}$$

$$\lim_{N \rightarrow \infty} \int d^D r r C_{stat}(r) = \infty \quad \Rightarrow \text{reached by coarsening}$$

INTERPOLATING BETWEEN ISING AND VOTER



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INTERPOLATING BETWEEN ISING AND VOTER

- $p = 1$ is voter
- $p = 2$ maps onto voter
- $p \geq 3$ Ising

Computational time $\sim pN$



CONCLUSIONS

- Space dependent interactions profoundly modify the coarsening process
- Much more remains to be done beyond determination of $L(t)$

- Thank you

